Another two-parameter excess Gibbs energy model was developed from an expansion of $(RTx_1x_2)/G^E$ instead of G^E/RTx_1x_2 . The end results are: $G^E = A_{12}^{/}A_{21}^{/}$

$$RTx_1x_2 = A_{12}'x_1 + A_{21}'x_2$$
 (12.16)

for the excess Gibbs energy and:

$$\ln \gamma_{1} = A_{12}^{/} \left(1 + \frac{A_{12}^{/} x_{1}}{A_{21}^{/} x_{2}} \right)^{-2}$$
(12.17a)
$$\ln \gamma_{2} = A_{21}^{/} \left(1 + \frac{A_{21}^{/} x_{2}}{A_{12}^{/} x_{1}} \right)^{-2}$$
(12.17b)

for the activity coefficients.

Note that: as
$$x_1 \rightarrow 0$$
, $\ln \gamma_1^{\infty} \rightarrow A'_{12}$

and as
$$x_2 \rightarrow 0$$
, $\ln \gamma_2^{\infty} \rightarrow A'_{21}$

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P5. Vapour liquid data for the system 1,4 dioxane(1)/ethylbenzene(2) at 85 °C are provided below. From these data obtain estimates of the van Laar coefficients (estimates based on smoothly drawn curves on the enclosed graph paper are sufficient). Estimate values of P-x₁-y₁ for x₁=0.5 based on these parameter values. Table 2: VLE data for the system 1,4 dioxane/ethylbenzene at 85 °C

P (kPa)	x ₁	У ₁
20.2	0	0
23.4	0.065	0.15
35.2	0.29	0.55
41.6	0.44	0.7
52.4	0.78	0.91
57.8	1	1

Local Composition Models

Unfortunately, the previous approach cannot be extended to systems of 3 or more components. For these cases, local composition models are used to represent multi-component systems.

- ➢ Wilson's Theory
- Non-Random-Two-Liquid Theory (NRTL)
- Universal Quasichemical Theory (Uniquac)

While more complex, these models have two advantages:

- > the model parameters are temperature dependent
- the activity coefficients of species in multi-component liquids can be calculated using information from binary data.



Wilson's Equations for Binary Solution Activity

A versatile and reasonably accurate model of excess Gibbs Energy was developed by Wilson in 1964. For a binary system, G^E is provided by:

$$\frac{G^{E}}{RT} = x_{1} \ln(x_{1} + x_{2}\Lambda_{12}) - x_{2} \ln(x_{2} + x_{1}\Lambda_{21})$$
(12.18)

where

$$\Lambda_{12} = \frac{V_2}{V_1} \exp\left[\frac{-a_{12}}{RT}\right] \qquad \Lambda_{21} = \frac{V_1}{V_2} \exp\left[\frac{-a_{21}}{RT}\right] \quad (12.24)$$

 V_i is the molar volume at T of the pure component *i*. a_{ij} is determined from experimental data.

The notation varies greatly between publications. This includes, $a_{12} = (\lambda_{12} - \lambda_{11}), a_{21} = (\lambda_{12} - \lambda_{22})$ that you will encounter in Holmes, M.J. and M.V. Winkle (1970) *Ind. Eng. Chem.* 62, 21-21. Wilson's Equations for Binary Solution Activity

Activity coefficients are derived from the excess Gibbs energy using the definition of a partial molar property:

$$RT \ln \gamma_{i} = \overline{G}_{i}^{E} = \frac{\partial n G^{E}}{\partial n_{i}} \bigg|_{T,P,n_{j}}$$

When applied to equation 11.16, we obtain:

$$\ln \gamma_1 = -\ln(x_1 + x_2 \Lambda_{12}) + x_2 \left(\frac{\Lambda_{12}}{x_1 + x_2 \Lambda_{12}} - \frac{\Lambda_{21}}{x_2 + x_1 \Lambda_{21}} \right) \quad (12.19a)$$

$$\ln \gamma_2 = -\ln(x_2 + x_1 \Lambda_{21}) - x_1 \left(\frac{\Lambda_{12}}{x_1 + x_2 \Lambda_{12}} - \frac{\Lambda_{21}}{x_2 + x_1 \Lambda_{21}} \right) \quad (12.19b)$$

Wilson's Equations for Multi-Component Mixtures

The strength of Wilson's approach resides in its ability to describe multi-component (3+) mixtures using binary data.

- Experimental data of the mixture of interest (ie. acetone, ethanol, benzene) is not required
- We only need data (or parameters) for acetone-ethanol, acetone-benzene and ethanol-benzene mixtures

The excess Gibbs energy for multicomponent mixtures is written:

$$\frac{G^E}{RT} = -\sum_i x_i \ln(\sum_j x_j \Lambda_{ij})$$
(12.22)

and the activity coefficients become:

$$\ln \gamma_{i} = 1 - \ln \sum_{i} x_{j} \Lambda_{ij} - \sum_{k} \frac{x_{k} \Lambda_{ki}}{\sum_{i} x_{j} \Lambda_{kj}}$$
(12.23)

where $\Lambda_{ij} = 1$ for i=j. Summations are over all species.

For three component systems, activity coefficients can be calculated from the following relationship:

$$\ln \gamma_{i} = 1 - \ln(x_{1}\Lambda_{i1} + x_{2}\Lambda_{i2} + x_{3}\Lambda_{i3}) - \frac{x_{1}\Lambda_{1i}}{x_{1} + x_{2}\Lambda_{12} + x_{3}\Lambda_{13}} - \frac{x_{2}\Lambda_{2i}}{x_{1}\Lambda_{21} + x_{2} + x_{3}\Lambda_{23}} - \frac{x_{3}\Lambda_{3i}}{x_{1}\Lambda_{31} + x_{2}\Lambda_{32} + x_{3}}$$

Model coefficients are defined as $(\Lambda_{ij} = 1 \text{ for } i=j)$:

$$\Lambda_{ij} = \frac{V_j}{V_i} \exp\left[\frac{-a_{ij}}{RT}\right]$$

Comparison of Liquid Solution Models

Activity coefficients of 2-methyl-2-butene + n-methylpyrollidone.

Comparison of experimental values with those obtained from several equations whose parameters are found from the infinite-dilution activity coefficients.

- (1) Experimental data.
- (2) Margules equation.
- (3) van Laar equation.
- (4) Scatchard-Hamer equation.
- (5) Wilson equation.



We now have the tools required to describe and calculate vapourliquid equilibrium conditions for even the most non-ideal systems.

1. Equilibrium Criteria:

In terms of chemical potential

$$\dot{\mu}_{i}^{v} = \mu_{i}^{l}$$

In terms of mixture fugacity $\hat{\mathbf{f}}_{i}^{\mathsf{v}} = \hat{\mathbf{f}}_{i}^{\mathsf{I}}$

2. Fugacity of a component in a non-ideal gas mixture:
$$\hat{f}_i^{\vee}(T,P,y_1,y_2,...,y_n) = y_i \hat{\phi}_i^{\vee}(T,P,y_1,y_2,...,y_n)P$$

3. Fugacity of a component in a non-ideal liquid mixture: $\hat{f}_{i}^{\dagger}(T,P,X_{1},X_{2},...,X_{n}) = X_{i}\gamma_{i}(T,P,X_{1},X_{2},...,X_{n})f_{i}^{\dagger}$

$$= \mathbf{X}_{i} \gamma_{i} (\mathsf{T}, \mathsf{P}, \mathbf{X}_{1}, \mathbf{X}_{2}, ..., \mathbf{X}_{n}) \phi_{i}^{\text{sat}} \mathsf{P}_{i}^{\text{sat}} \exp \left[\frac{\mathsf{V}_{i} (\mathsf{P} - \mathsf{P}_{i}^{\text{sat}})}{\mathsf{RT}} \right]$$
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γ, ϕ Formulation of VLE Problems

To this point, Raoult's Law was only description we had for VLE behaviour:

$$y_i P = x_i P_i^{\text{sat}}$$

We know that calculations based on Raoult's Law do not predict actual phase behaviour due to over-simplifying assumptions.

<u>Accurate</u> treatment of non-ideal phase equilibrium requires the use of mixture fugacities. At equilibrium, the fugacity of each component is the same in all phases. Therefore,

or,

$$\begin{aligned} \hat{f}_{i}^{v} &= \hat{f}_{i}^{I} \\ y_{i} \hat{\varphi}_{i}^{v} P &= x_{i} \gamma_{i} \varphi_{i}^{sat} P_{i}^{sat} exp \Bigg[\frac{V_{i} (P - P_{i}^{sat})}{RT} \Bigg] \end{aligned}$$

determines the VLE behaviour of non-ideal systems where Raoult's Law fails.

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Non-Ideal VLE to Moderate Pressures

A simpler expression for non-ideal VLE is created upon defining a lumped parameter, Φ .

$$\Phi_{i} = \frac{\hat{\phi}_{i}^{v}}{\phi_{i}^{sat}} \exp\left[\frac{-V_{i}^{l}(P - P_{i}^{sat})}{RT}\right]$$
$$\approx \frac{\hat{\phi}_{i}^{v}}{\phi_{i}^{sat}}$$

The final expression becomes,

$$y_i \Phi_i P = x_i \gamma_i P_i^{sat}$$
 (i = 1,2,3,...,N) 14.1

To perform VLE calculations we therefore require vapour pressure data (P_i^{sat}), vapour mixture and pure component fugacity correlations (Φ_i) and liquid phase activity coefficients (γ_i).

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Sources of Data:

1. Vapour pressure: Antoine's Equation (or similar correlations)

$$lnP_{i}^{sat} = A_{i} - \frac{B_{i}}{T + C_{i}}$$
 14.3

2. Vapour Fugacity Coefficients: Viral EOS (or others)

$$\Phi_{i} = exp\left[\frac{\mathsf{B}_{ii}(\mathsf{P}-\mathsf{P}^{sat}_{i}) + 0.5\mathsf{P}\sum_{j} y_{j}y_{k}(2\delta_{ji} - \delta_{jk})}{\mathsf{RT}}\right] \qquad 14.6$$

- 3. Liquid Activity Coefficients
 - Binary Systems Margules, van Laar, Wilson, NRTL, Uniquac
 - > Ternary (or higher) Systems Wilson, NRTL, Uniquac